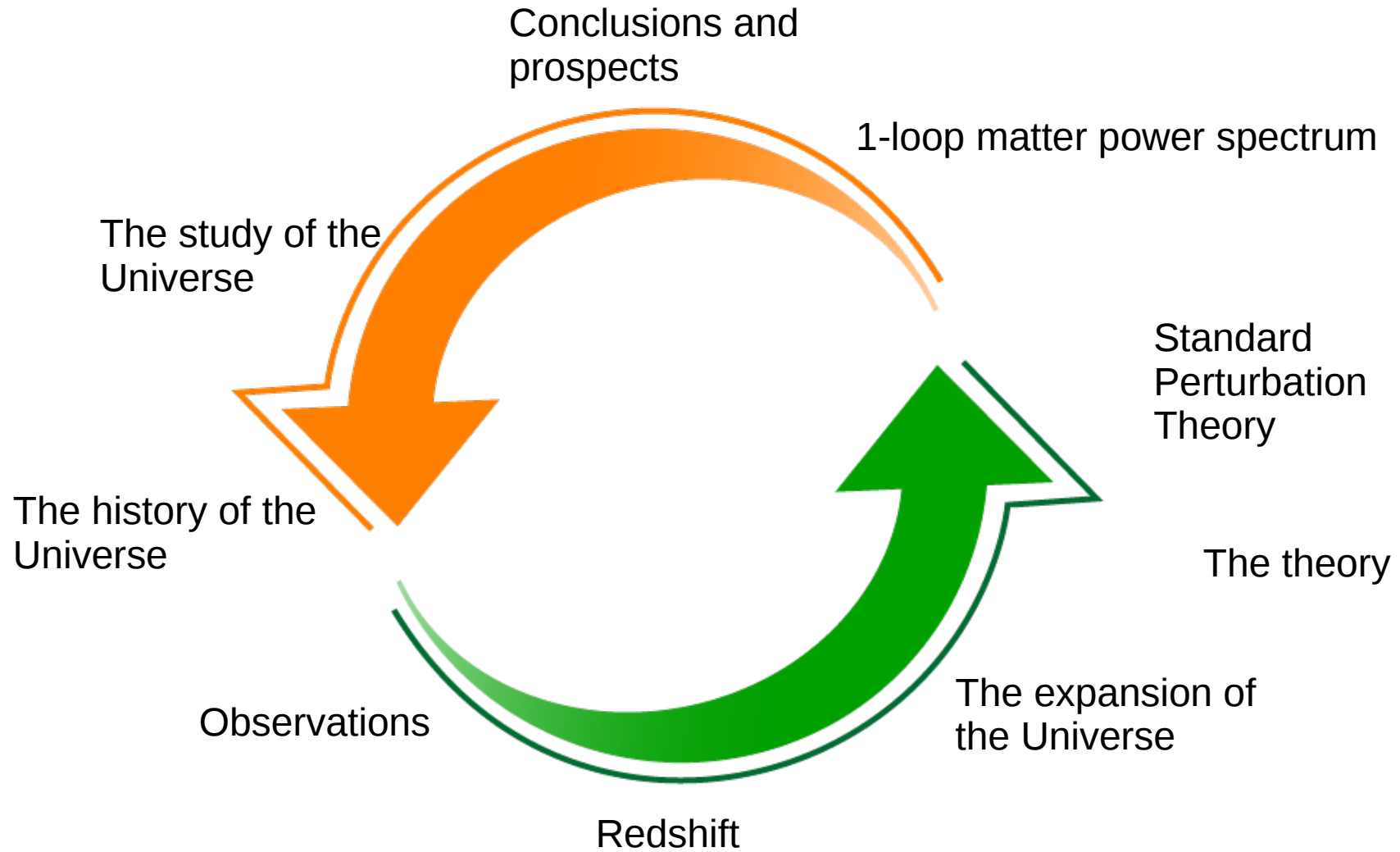


THE UNIVERSE in a COCONut-shell

**STANDARD THEORY
OF COSMOLOGICAL
PERTURBATIONS**

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University of Sussex

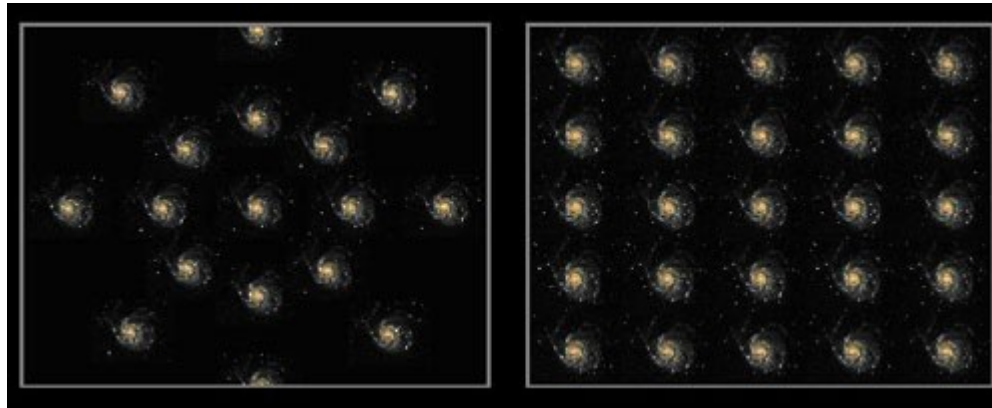
Outline



The study of the Universe

And the distribution of matter throughout the Universe

Isotropic
(rotational
invariance)



Homogeneous
(translational
invariance)

Cosmological
principle

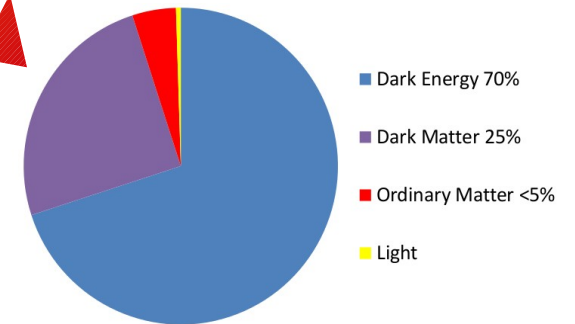
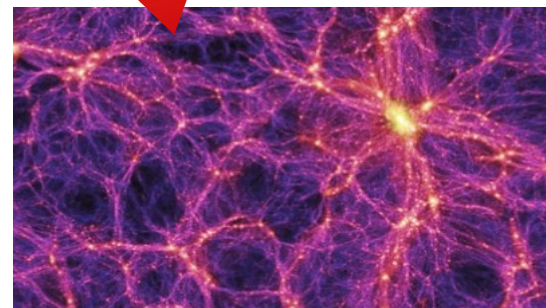
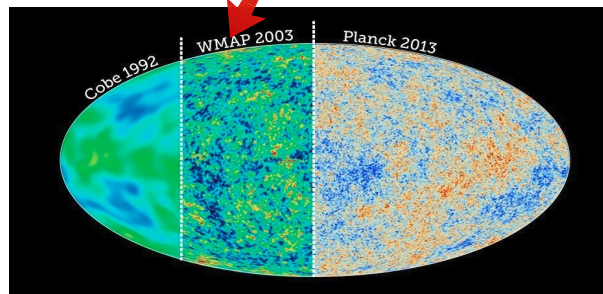
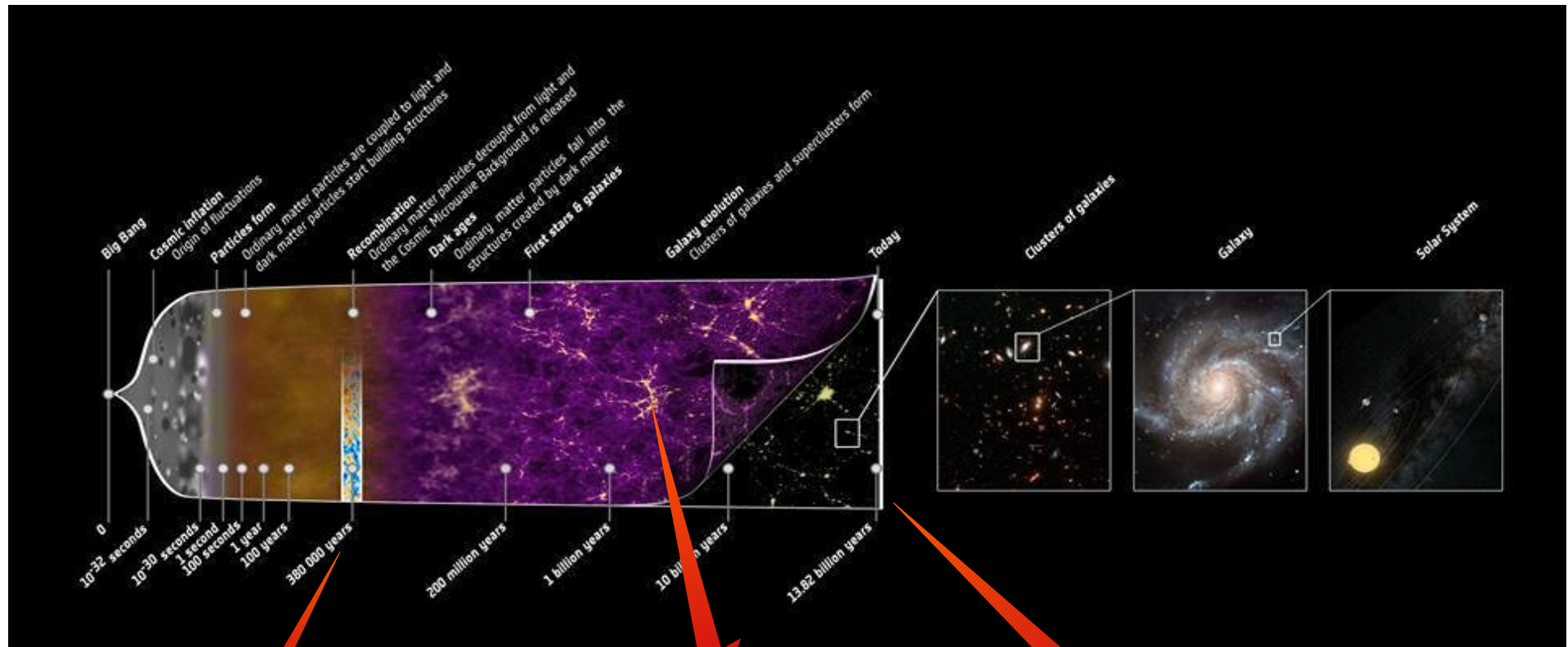
Real Universe

**Small asymmetries
and fluctuations**
risen from the beginning
of time



STRUCTURES
today!

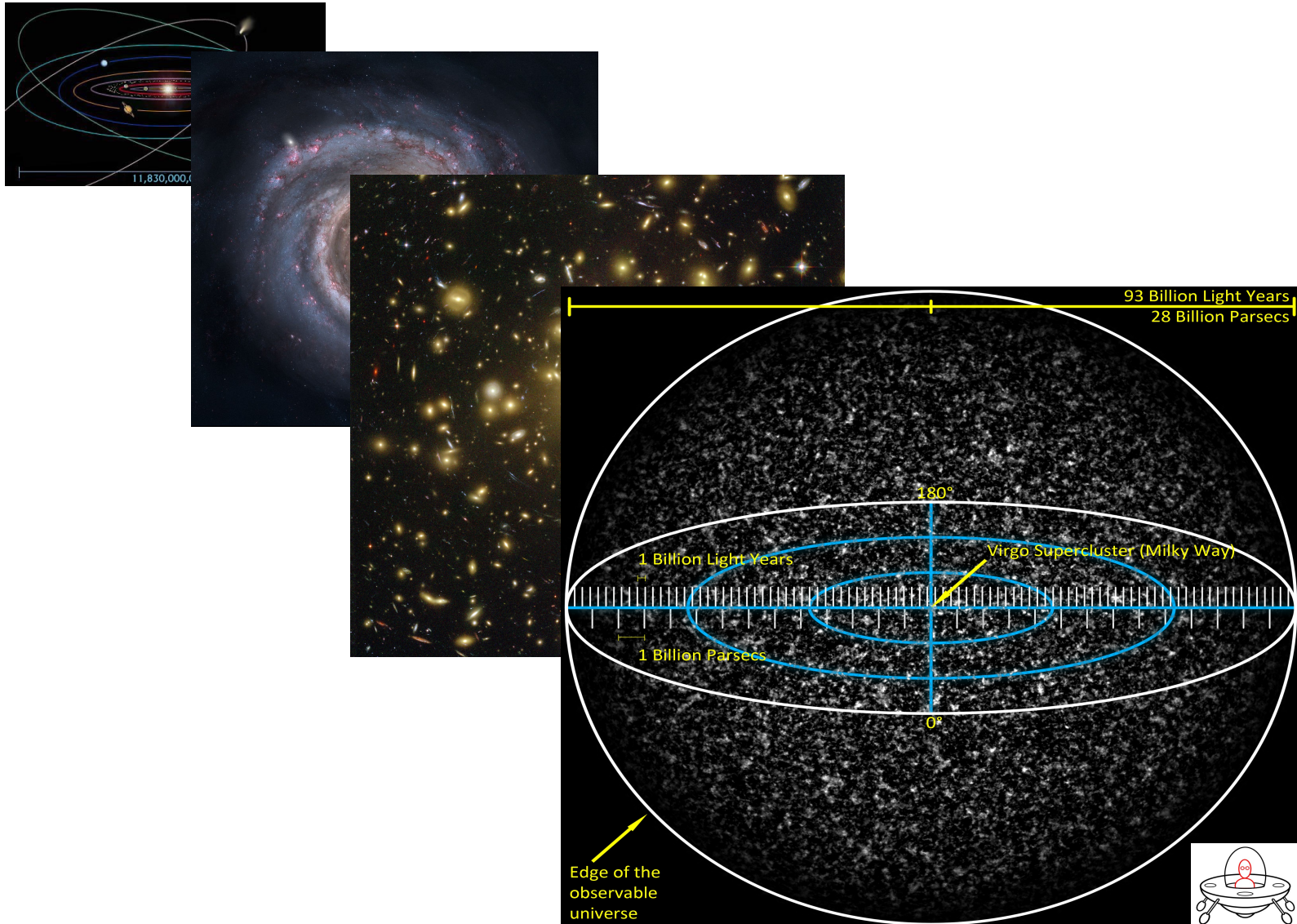
The history of our Universe



I. OBSERVATIONS

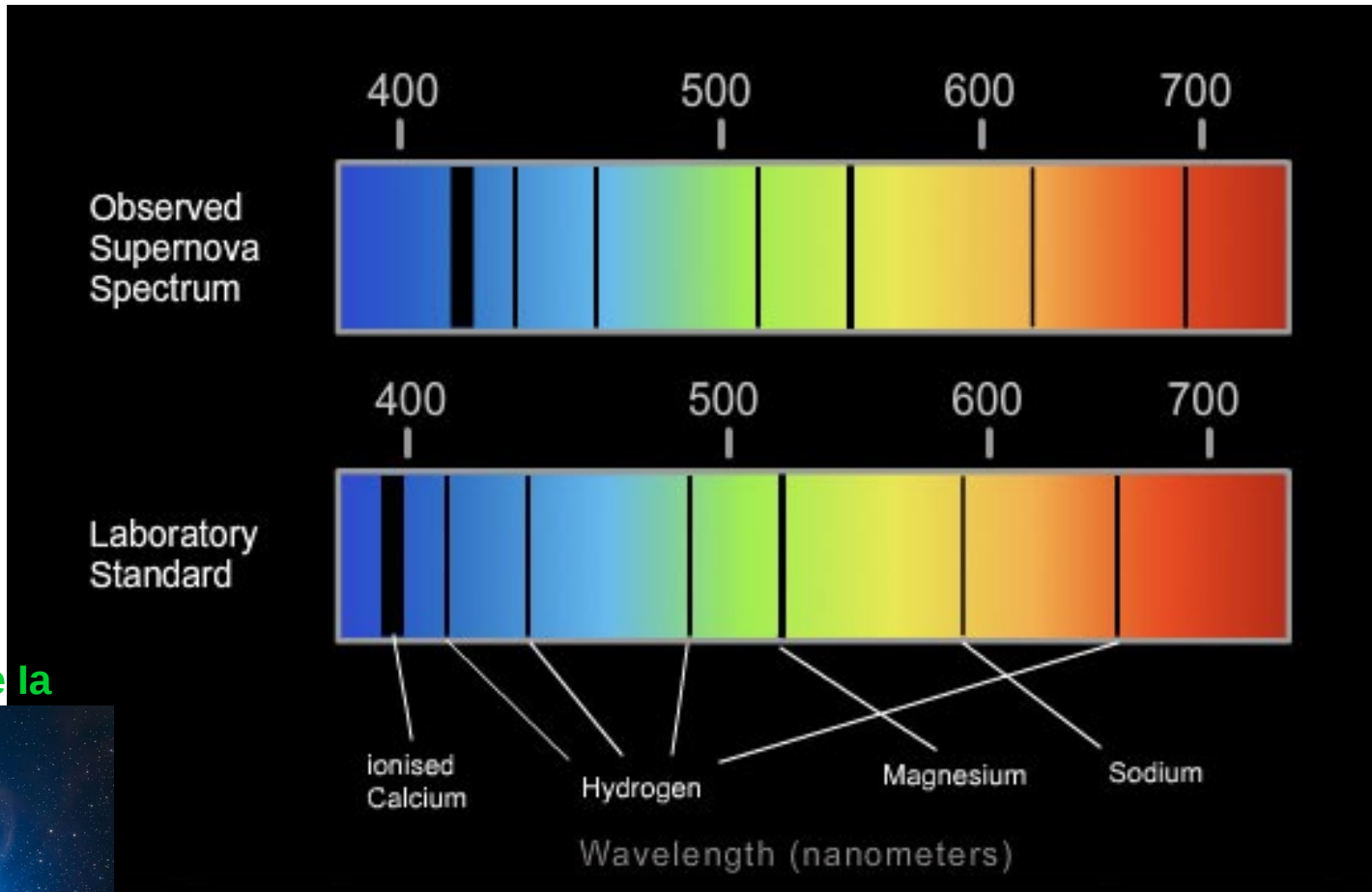


Observations

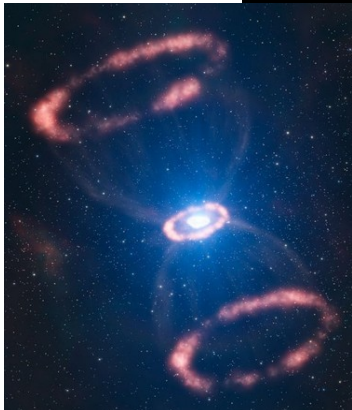


Redshift

$$z = \frac{\lambda_o - \lambda_e}{\lambda_e}$$



Supernovae Ia



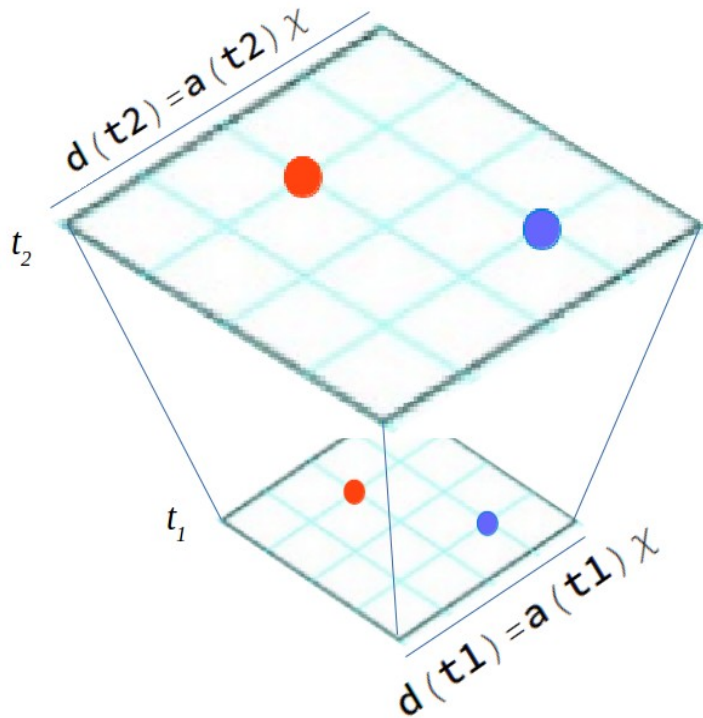
Redshift

Distances

Time

Z=0 means "present time"

The expansion of the Universe



Fundamental observers:

their coordinates do not change with the expansion of the universe.

1. Co-moving coordinate system (grid)
Co-moving distance, χ .

2. Physical distances, $d(t)$, change with the **scale factor**, $a(t)$, as the universe expands.

Expansion rate or **Hubble parameter**: $H(t) = \frac{\dot{a}}{a}$

Redshift – scale factor $a(t) = \frac{1}{1+z}$

Hubble's law

Hubble flow (recession) **velocity**:

$$\frac{\partial d(t)}{\partial t} \equiv v = H d(t)$$

Peculiar velocity: (anisotropy)

$$v_{pec} = a(t) \frac{\partial \chi}{\partial t} \rightarrow v = H d(t) + v_{pec}$$

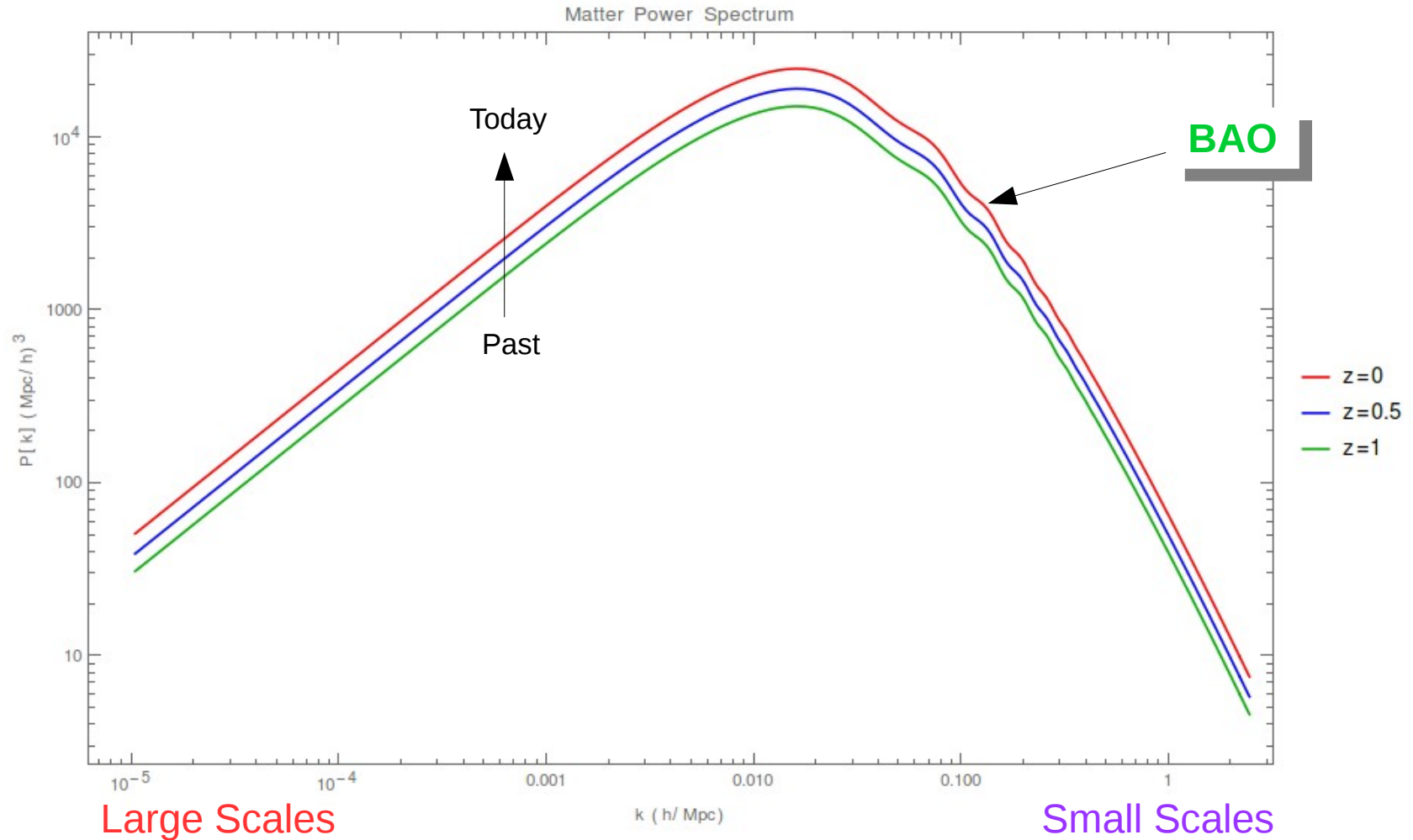


**Relation
redshift - distances**

Power Spectrum

WHAT TO DO WITH ALL
THIS INFORMATION?

Power Spectrum



Planck 2015 cosmology

CAMB [A. Lewis]

II. STANDARD PERTURBATION THEORY

Description of observations

General Relativity

How to measure distances?

Flat, homogeneous and isotropic universe

$$ds^2 = -e^{2\Psi} dt^2 + a^2 e^{2\Phi} dx^2$$

Variational principle



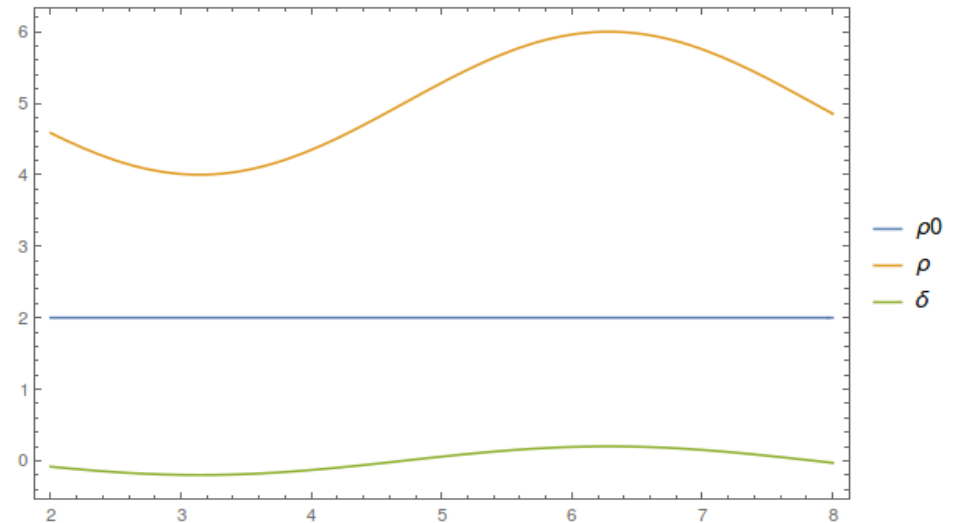
Fluid equations + Poisson constraint

energy density ρ and pressure P

Dark Matter (non relativistic)

$$P = 0$$
$$\rho = \rho_0 + \delta \rho$$
$$\delta = \frac{\delta \rho}{\rho}$$

Density contrast



Fluid equations

Density conservation

$$\dot{\delta} + \nabla \cdot ((1 + \delta) \underline{v}) = 0 \quad [1]$$

Newton's law

$$\dot{\underline{v}} + (\underline{v} \cdot \nabla) \underline{v} + 2H \underline{v} - \frac{1}{a^2} \nabla \Phi = 0 \quad [2]$$

Poisson's equation

$$\frac{1}{a^2} \nabla^2 \Phi = -\frac{3H^2}{2} \Omega_m \delta \quad [3]$$

Cooking up:

$$\Omega_m = \rho_m / 3H^2 M_p^2$$

1. Re-arrange terms to get a single second-order equation for delta.
2. Fourier space.

3. Time derivatives $\dot{\square} \equiv \frac{\partial}{\partial t}$ \rightarrow redshift derivatives $\square' \equiv \frac{\partial}{\partial z} = -\frac{H}{a} \frac{\partial}{\partial t}$

4. Assumptions

Perfect fluid behaviour,
non-relativistic limit,
radial inflow and
negligible vorticity.



The matter density equation

$$\epsilon = -\frac{\dot{H}}{H^2}$$

$$\delta_k'' - \frac{1-\epsilon}{1+z} \delta_k' - \frac{3}{2} \frac{\Omega_m}{(1+z)^2} \delta_k =$$

Linear equation

Non-linear contribution

$$\begin{aligned} & - \int \frac{d^3 \mathbf{q} d^3 \mathbf{s}}{(2\pi)^6} (2\pi)^3 \delta(\mathbf{k} - \mathbf{q} - \mathbf{s}) S_2(\mathbf{q}, \mathbf{s}) \\ & - \int \frac{d^3 \mathbf{q} d^3 \mathbf{s}}{(2\pi)^6} (2\pi)^3 \delta(\mathbf{k} - \mathbf{q} - \mathbf{s}) \int \frac{d^3 \mathbf{\tau} d^3 \mathbf{\sigma}}{(2\pi)^6} (2\pi)^3 \delta(\mathbf{s} - \mathbf{\tau} - \mathbf{\sigma}) S_3(\mathbf{q}, \mathbf{s}, \mathbf{\tau}, \mathbf{\sigma}) \\ & + O(\delta^4) \end{aligned} \quad [4]$$

Green's method

Solving the equation... via expansion in delta $O(\delta^3)$ (~1-loop)

Linear solution

$$\delta^{(1)}(k, z) = D(z) \tilde{\delta}_k$$

Initial conditions in the past
within matter-dominated era

$$\tilde{\delta}_k = \delta_k(\tilde{z}=50)$$

$$D'' - \frac{1-\epsilon}{1+z} D' - \frac{3}{2} \frac{\Omega_m}{(1+z)^2} D = 0$$

Growth function

Second-order

Third-order solution

Insert	in	To get
$\delta_k^{(1)}$	$S_2(\mathbf{q}, \mathbf{s})$	$\delta^{(2)}(k, z)$
$\delta_k^{(1)}, \delta_k^{(2)}$	$S_3(\mathbf{q}, \mathbf{s}, \mathbf{\tau}, \mathbf{\sigma})$	$\delta^{(3)}(k, z)$

$$\delta = \delta^{(1)} + \delta^{(2)} + \delta^{(3)}$$

2-Point Correlation Function

$$\langle \delta^{(n)}(k_1, z) \delta^{(n)}(k_2, z) \rangle = (2\pi)^3 \delta(\underline{k}_1 + \underline{k}_2) P_{nn}(k, z)$$

Tree level

$$P_{11}(k, z) = D(z)^2 \tilde{P}(k) \quad [5]$$

1 loop corrections

$$P_{22}(k, z) = D_A^2(z) P_{AA}(k, z) + D_A(z) D_B(z) P_{AB}(k, z) + D_B^2(z) P_{BB}(k, z) \quad [6]$$

$$P_{13}(k, z) = D(z) \tilde{P}(k) [(D_D(z) - D_J(z)) P_D(k, z) + D_E(z) P_E(k, z)] \\ + D(z) \tilde{P}(k) [(D_F(z) + D_J(z)) P_F(k, z) + D_G(z) P_G(k, z)] \\ + D(z) \tilde{P}(k) \left[\frac{D_J(z)}{2} (P_{J2}(k, z) - 2P_{J1}(k, z)) \right] \quad [7]$$

$\tilde{P}(k) \equiv P(k, \tilde{z}=50)$

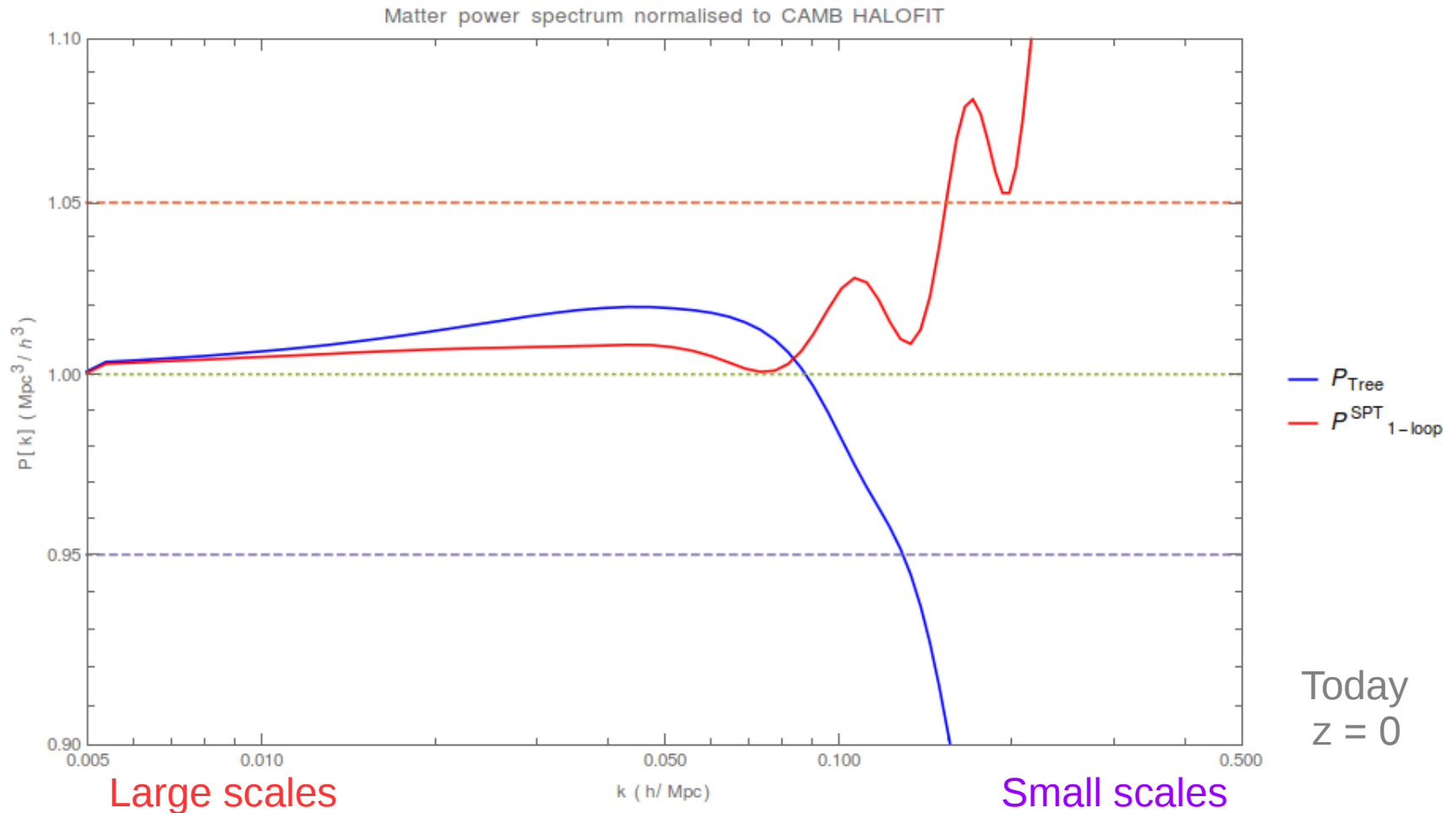
Initial power spectrum
by CAMB

$D_{IJ} \forall I, J = A, B$ AND $D_l \forall l = D, E, F, G, J1, J2$ BEING GROWTH FUNCTIONS

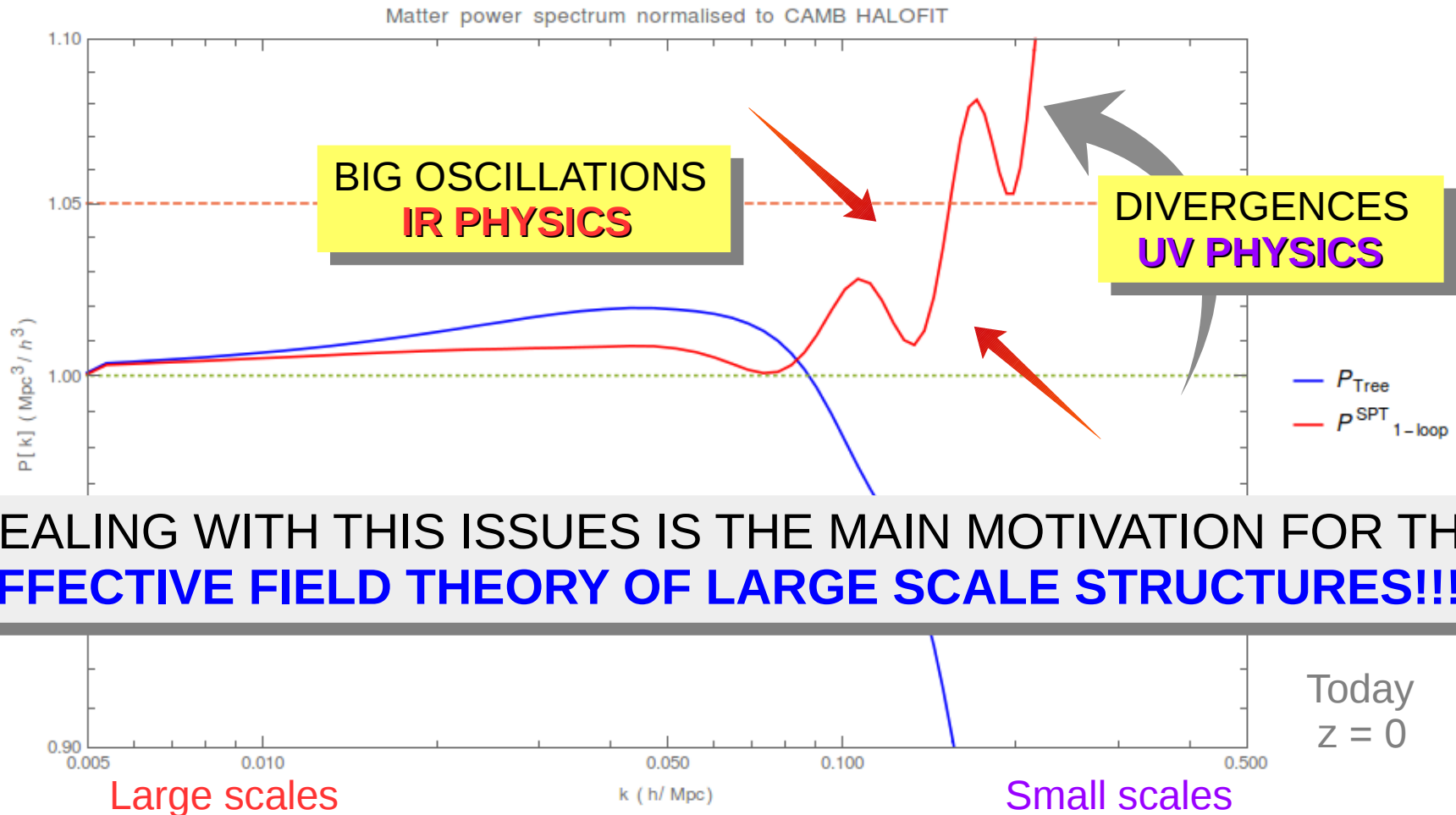
$P_{IJ} \forall I, J = A, B$ AND $P_l \forall l = D, E, F, G, J1, J2$ BEING LOOP INTEGRALS

$$P_{1-loop}^{SPT}(k, z) = P_{11}(k, z) + P_{13}(k, z) + P_{22}(k, z) \quad [8]$$

1-loop matter power spectrum



1-loop matter power spectrum



Conclusions and Prospects

- The 1-loop matter power spectrum involves weighted integrals over the initial power spectrum with certain weighting functions.
- Similarly to Quantum Field Theory, the influence of the amplitude of fluctuations can be very large or even unbounded in the ultra-violet regime.
- Unreliable predictions: ignorance of an ultimate high-energy theory and troubles to accurately model details of non-linear halo, galaxy formation, gas dynamics...
- We need a re-normalization procedure to account for the effects of ultra-violet physics.
- IR-resummation?
- What about Redshift Space Distortions?