LIMBER **&**

THE JOURNEY OF THE Sky PROJECT

Lucia F. de la Bella 09/02/2021



https://orcid.org/0000-0002-1064-3400 https://howtoreachthecosmos.jimdofree.com

Also known as

Lucia F. de la Bella

Lucía Fonseca Lucía Fonseca de la Bella



Who I am...

MANCHESTER 1824

The University of Manchester

PDRA



Weak Lensing Unequal-time correlators



2018-201

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THE JOURNEY OF LIMBER

L.F. de la Bella, N. Tessore and S. Bridle (arXiv 2011.06185) N. Tessore and L. F. de la Bella (in prep)

> Python packages UNEQUALPY CORFU







1. Cosmic Lensing



Correlations between fields

- Same time slice: *equal-time correlators* •
- Different time slices: *unequal-time correlators* Examples of fields:

Matter, convergence, cosmic shear

Note: lookback time = redshift = co-moving distance

Angular correlation function

$$w(\theta) = \iint_{0}^{\infty} dx_1 \, dx_2 \, f_1(x_1) f_2(x_2) \, \xi(r_{12}; t_1, t_2)$$

FILTERS

UNEQUAL-TIME **CORRELATION FUNCTION**

Angular power spectrum

 $C(\ell) = \int_0^\infty \frac{dk}{k^2} \iint_0^\infty dx_1 \, dx_2 \, f_1(x_1) f_2(x_2) j_{\ell}(kx_1) j_{\ell}(kx_2) \, P(k; t_1, t_2)$

Bessel Functions

UNEQUAL-TIME POWER SPECTRUM



1. Cosmic Lensing



Correlations between fields

- Same time slice: *equal-time correlators* •
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$$w(\theta) = \iint_{0}^{\infty} dx_{1} dx_{2} f_{1}(x_{1}) f_{2}(x_{2}) \xi(r_{12}; t_{1}, t_{2})$$
$$C(\ell) = \int_{0}^{\infty} \frac{dk}{k^{2}} \iint_{0}^{\infty} dx_{1} dx_{2} f_{1}(x_{1}) f_{2}(x_{2}) j_{\ell}(kx_{1}) j_{\ell}(kx_{2}) P(k; t_{1}, t_{2})$$



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2. The Journey

- Filters
- **Bessel functions**
- Unequal-time correlators •

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$$w(\theta) = \iint_{0}^{\infty} dx_{1} dx_{2} f_{1}(x_{1}) f_{2}(x_{2}) \xi(r_{12}; t_{1}, t_{2})$$
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- i) Filters
- ii) Mid point

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2. The Journey

Assumptions:

- Smooth filters i)
- Correlation falls off fast ii)
- iii) Small angle separation



$w(\theta) \approx \int_0^\infty d\bar{x} \int_{-2\bar{x}}^{2\bar{x}} dR f_1(\bar{x}) f_2(\bar{x}) \,\xi(r_{12};\bar{t})$





$$w(\theta) = \iint_{0}^{\infty} dx_{1} dx_{2} f_{1}(x_{1}) f_{2}(x_{2}) \xi(r_{12}; t_{1}, t_{2})$$
$$C(\ell) = \int_{0}^{\infty} \frac{dk}{k^{2}} \iint_{0}^{\infty} dx_{1} dx_{2} f_{1}(x_{1}) f_{2}(x_{2}) j_{\ell}(kx_{1}) j_{\ell}(kx_{2}) P(k; t_{1}, t_{2})$$

$$w(\theta) \approx \int_0^\infty d\bar{x} \int_{-2\bar{x}}^{2\bar{x}} dR f_1(\bar{x}) f_2(\bar{x}) \xi(r_{12};\bar{t})$$

i) Filters

- i) Sphere
- ii) Mid point
- ii) Discrete case

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PEEBLES



2. The Journey

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$$w(\theta) = \iint_{0}^{\infty} dx_{1} dx_{2} f_{1}(x_{1}) f_{2}(x_{2}) \xi(r_{12}; t_{1}, t_{2})$$
$$C(\ell) = \int_{0}^{\infty} \frac{dk}{k^{2}} \iint_{0}^{\infty} dx_{1} dx_{2} f_{1}(x_{1}) f_{2}(x_{2}) j_{\ell}(kx_{1}) j_{\ell}(kx_{2}) P(k; t_{1}, t_{2})$$

$$w(\theta) \approx \int_0^\infty d\bar{x} \int_{-2\bar{x}}^{2\bar{x}} dR f_1(\bar{x}) f_2(\bar{x}) \xi(r_{12};\bar{t})$$

- i) Filters
- ii) Mid point
- i) Sphere
- ii) Discrete case
- i) Flat sky ii) Curved universe
- iii) Dirac delta

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2. The Journey

$$C(\ell) = \int d^2 \vec{\theta} w(\theta) e^{-i\vec{\ell} \cdot \vec{\theta}}$$

$$\approx \int dx \frac{f_1(x) f_2(x)}{x^2} P(\ell/x; t)$$
Fourier space
Delta functions





$$w[\theta] = \iint_{0}^{\infty} dx_{1} dx_{2} f_{1}(x_{1}) f_{2}(x_{2}) \xi(r_{12}; t_{1}, t_{2})$$
$$C(\ell) = \int_{0}^{\infty} \frac{dk}{k^{2}} \iint_{0}^{\infty} dx_{1} dx_{2} f_{1}(x_{1}) f_{2}(x_{2}) j_{\ell}(kx_{1}) j_{\ell}(kx_{2}) P(k; t_{1}, t_{2})$$



9

0,0

10,5

- Thin-layer approx.
- iii) Limber's inaccurate

SIMON

2. The Journey

Accuracy of Limber:

- Small angles: Limber
- Large angles: Thin-layer approximation
- Compares Limber's vs "exact"
- Applies small-angle to "exact"

i) Small and large angles





$$w(\theta) = \iint_{0}^{\infty} dx_{1} dx_{2} f_{1}(x_{1}) f_{2}(x_{2}) \xi(r_{12}; t_{1}, t_{2})$$
$$C(\ell) = \int_{0}^{\infty} \frac{dk}{k^{2}} \iint_{0}^{\infty} dx_{1} dx_{2} f_{1}(x_{1}) f_{2}(x_{2}) \frac{j_{\ell}(kx_{1}) j_{\ell}(kx_{2})}{j_{\ell}(kx_{1}) j_{\ell}(kx_{2})} P(k; t_{1}, t_{2})$$



$$C(\ell) \approx \int dx \frac{f_1(x)f_2(x)}{x^2} P(\ell/x;t)$$



2. The Journey

Dirac delta version:

$$j_{\ell}(kx) \approx \sqrt{\frac{\pi}{2\nu}} \delta_D(kx - \nu)$$

Series expansion in $1/\nu$

$$\nu = \ell + 1/2$$



$$w(\theta) = \iint_{0}^{\infty} dx_{1} dx_{2} f_{1}(x_{1}) f_{2}(x_{2}) \xi(r_{12}; t_{1}, t_{2})$$
$$C(\ell) = \int_{0}^{\infty} \frac{dk}{k^{2}} \iint_{0}^{\infty} dx_{1} dx_{2} f_{1}(x_{1}) f_{2}(x_{2}) \frac{j_{\ell}(kx_{1}) j_{\ell}(kx_{2})}{j_{\ell}(kx_{1}) j_{\ell}(kx_{2})} P(k; t_{1}, t_{2})$$

$$w(\theta) \approx \int_0^\infty d\bar{x} \int_{-2\bar{x}}^{2\bar{x}} dR f_1(\bar{x}) f_2(\bar{x}) \xi(r_{12};\bar{t})$$

i) Filters

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ii) Mid point

i) Sphere ii) Discrete case

PEEBLES

$$C(\ell) \approx \int dx \frac{f_1(x)f_2(x)}{x^2} P(\ell/x;t)$$

i) Flat sky

- i) Small and larg
- ii) Curved universe ii) Thin-layer app
- iii) Dirac delta

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2. The Journey

Geometric approximation (large scales)

$$P(k; t_1, t_2) \approx \sqrt{P(k, t_1)P(k, t_2)} \qquad C(\ell) \approx \frac{\pi}{2\nu} \int \frac{dk}{k^2} f_1(\nu/k) f_2(\nu/k) P(k, \nu/k)$$

$$i) \quad Small \text{ and large angles}$$

$$i) \quad Post-Limber$$

$$i) \quad Post-Limber$$

$$i) \quad Post-Limber$$

$$i) \quad Divergence \quad small \quad \ell \quad ii) \quad Dirac \quad delta$$

$$ii) \quad Divergence \quad small \quad \ell \quad iii) \quad Dirac \quad delta$$

$$ii) \quad Valid \quad Iarge \quad \ell$$

$$iv) \quad Valid \quad small \quad k$$

$$SIMON \qquad LOVERDE \qquad LEMOS$$

$$ET \quad AL \qquad ET \quad AL$$

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k)

$$w(\theta) = \iint_{0}^{\infty} dx_1 \, dx_2 \, f_1(x_1) f_2(x_2) \, \xi(r_{12}; t_1, t_2)$$
$$C(\ell) = \iint_{0}^{\infty} \frac{dk}{k^2} \iint_{0}^{\infty} dx_1 \, dx_2 \, f_1(x_1) f_2(x_2) \, j_{\ell}(kx_1) \, j_{\ell}(kx_2) \, P(k; t_1, t_2)$$

$$w(\theta) \approx \int_0^\infty d\bar{x} \int_{-2\bar{x}}^{2\bar{x}} dR f_1(\bar{x}) f_2(\bar{x}) \xi(r_{12};\bar{t})$$

- ii) Mid point

$$C(\ell) \approx \int dx \frac{f_1(x)f_2(x)}{x^2} P(\ell/x;t)$$



2. The Journey





$$w(\theta) = \iint_{0}^{\infty} dx_{1} dx_{2} f_{1}(x_{1}) f_{2}(x_{2}) \xi(r_{12}; t_{1}, t_{2})$$
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- ii) Mid point

$$C(\ell) \approx \int dx \frac{f_1(x)f_2(x)}{x^2} P(\ell/x;t)$$



2. The Journey





N. Tessore, L. F. de la Bella (in prep) Python package **CORFU**



3. All-angle approach

$$w(\theta) = \iint_{0}^{\infty} dx_{1} dx_{1}$$
$$C(\ell) = \int_{0}^{\infty} \frac{dk}{k^{2}} \iint_{0}^{\infty} dx_{1} dx_{2}$$

- Need accurate unequal-time power spectrum
- Deal with non-linear physics (one-loop, EFT)
- Impact on weak lensing. How?
 - Midpoint approximation \star
 - Compare with geometric approx.
- Analyse regimes of validity
- Drop approximations
- Python package **UNEQUALPY**

- $dx_2 f_1(x_1) f_2(x_2) \,\xi(r_{12};t_1,t_2)$
- $f_1(x_1)f_2(x_2)j_{\ell}(kx_1)j_{\ell}(kx_2) P(k;t_1,t_2)$
 - Need exact calculation at all angular scales
 - Python package **CORFU**
 - FFTLog methods
 - Inverse Fourier transform

$$P(k; t_1, t_2) \longrightarrow \xi(r; t_1, t_2)$$

• Legendre polynomials

$$w(\theta) \longrightarrow C(\mathcal{E})$$



L. F. de la Bella, N. Tessore and S. L. Bridle (arXiv 2011.06185) Python package **UNEQUALPY**





4. Unequal-time Correlators

2point functions

Homogeneous Isotropic

EQUAL-TIME POWER SPECTRUM

Same time slice

UNEQUAL-TIME POWER SPECTRUM

Different time slices

LINEAR THEORY $P(k, z) = D(z)^2 P_{11}(k)$

$$P(k, z_1, z_2) = D(z_1)D(z_2)P_{11}(k)$$

(a)
$$z_1 = z_2$$

$$P(k, z, z) = P(k, z)$$

$$(2\pi)^3 \delta_D(\mathbf{k}_1 + \mathbf{k}_2) P(k, z) := \langle \delta(\mathbf{k}_1, z) \delta^*(\mathbf{k}_2, z) \rangle$$

 $(2\pi)^3 \delta_D(\mathbf{k}_1 + \mathbf{k}_2) P(k, z_1, z_2) := \langle \delta(\mathbf{k}_1, z_1) \delta^*(\mathbf{k}_2, z_2) \rangle$

Non-linear

ONE-LOOP THEORY

 $P(k, z_1, z_2) = P_{11}(k, z_1, z_2) + P_{22}(k, z_1, z_2) + P_{13}(k, z_1, z_2)$

- i) STANDARD PERTURBATION THEORY
- **EFFECTIVE FIELD THEORY** ii)

NL physics \rightarrow counterterms

$$P_{EFT} = P_{SPT} - c^2 k^2 P_{11}$$



- This work shows unequal-time EFT breaks
- New idea: the midpoint approximation!







Non-linear

*
$$P(k; t_1, t_2) \approx \sqrt{P(k, t_1)P(k, t_2)}$$

4.1. Time and non-linear effects



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Non-linear

*
$$P(k; t_1, t_2) \approx \sqrt{P(k, t_1)P(k, t_2)}$$

L)

4.1. Time and non-linear effects



$$\frac{1}{P^2} = \frac{1}{P_{theory}(k; z_1, z_2)^2}$$

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4.2. Midpoint approximation





4.2. Midpoint approximation





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Exact
$$C_{ab}^{(i,j)}(\ell) = \frac{2}{\pi} \int_0^\infty dk k^2 \iint_0^\infty dx$$

 $C_{ab}^{(i,j)}(\ell) \approx \frac{1}{\nu} \int_0^\infty dk k^2 f_a^i(\nu/k) f_b^j(\nu/k) P(k;\nu/k)$ Limber

Geometric
$$C_{ab}^{(i,j)}(\ell) \approx \frac{2}{\pi} \int_0^\infty dk k^2 \int_0^\infty dx_1 f_a^i(x_1) j_\ell(kx_1) \sqrt{P(k;x_1)} \int_0^\infty dx_2 f_b^j(x_2) j_\ell(kx_2) \sqrt{P(k;x_2)} dx_2 f_b^j(x_2) j_\ell(kx_2) \sqrt{P(k;x_2)} dx_2 f_b^j(x_2) dx_2 f_b^j($$

Midpoint $C_{ab}^{(i,j)}(\ell) \approx \frac{2}{\pi} \int_0^\infty dk k^2 \iint_0^\infty dx_1$

5. Results

$x_1 dx_2 f_a^i(x_1) f_b^j(x_2) j_{\ell}(kx_1) j_{\ell}(kx_2) P(k; t_1, t_2)$

CORFU $P(k; t_1, t_2) \longrightarrow \xi(r; t_1, t_2)$

$$w(\theta) \longrightarrow C(\mathcal{E})$$

$${}_{1} dx_{2} f_{a}^{i}(x_{1}) f_{b}^{j}(x_{2}) j_{\ell}(kx_{1}) j_{\ell}(kx_{2}) P\left(k; \frac{x_{1} + x_{2}}{2}\right)$$















5. Results

DES-Y1 DATA



6. Summary

Summary

Coming next!

- Angular correlations functions are very hard to compute
- Limber is the most widely used approximation
- List of issues: accuracy and validity of approximations
- Need for unequal-time correlators and all-angle computations
- Unequal-time EFT does not improve the prediction
- Midpoint approximation better on non-linear scales

• Numeric paper (in prep)

- All angular scale computations
- Equal and unequal-time correlators
- Phyton package CORFU

Science paper (arXiv 2011.06185)

- Unequal-time matter power spectrum at one-loop Phyton package UNEQUALPY
- Analysis of all approximations and validity regimes
- **Midpoint and Geometric** best to mimic unequal-time features!
- Beyond Limber relevant for galaxy clustering and galaxy-galaxy lensing.





Image by Simon Birrer

